



**POLITECNICO**  
MILANO 1863

# **6th Two-Day Meeting** **on Propulsion Simulations** **Using OpenFOAM Technology**



## **Modelling hydrogen injection for internal combustion engines using WENO schemes in the OpenFOAM framework**

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**Politecnico di Bari**

# Outlook

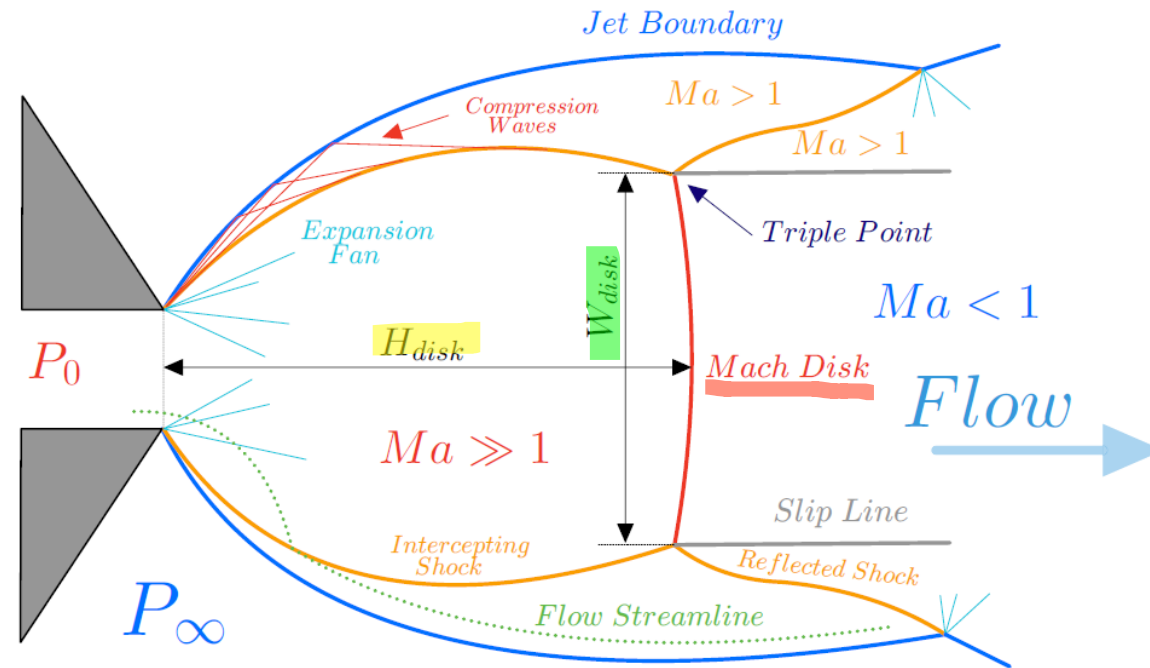
- Introduction
- Numerical method
- Simulation set-up
- Grid refinement study
- Results
- Conclusions
- Future work

# Introduction

- Hydrogen can be used in internal combustion engines to zero the carbon-based tailpipe exhaust emissions
- In-cylinder direct injection preferable to **avoid** preignition and backfire as well as to **obtain** high volumetric efficiency
- Hydrogen's low density requires **high injection pressures** for fast hydrogen penetration and sufficient in-cylinder mixing
- **Chocked flow conditions** during the injection which result in the formation of turbulent under-expanded jets
- Understanding of the underexpansion process and turbulent mixing just after the nozzle exit necessary for designing of an efficient hydrogen injection system and strategies

# Introduction

- Underexpanded jets of hydrogen in air for different Nozzle Pressure Ratios ( $\text{NPR} = P_0/P_\infty$ ) investigated
- Schematic of near-nozzle structure of under-expanded jets (picture is based on the visualization presented by Crist et al. [1])



- Very complex flow structures need high accurate scheme with low dissipation and good shock-capturing capabilities. Accurate prediction of H<sub>2</sub>-air mixing strongly depends on the description of the near-nozzle flow.

# Governing equations

The Favre-Reynolds averaged Navier-Stokes equations have been solved with k- $\omega$  SST turbulence model

$$\begin{aligned}\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j}{\partial x_j} &= 0, \\ \frac{\partial \bar{\rho} \tilde{u}_i}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{u}_i}{\partial x_j} &= - \frac{\partial \bar{p}}{\partial x_i} \\ &\quad - \frac{\partial}{\partial x_j} \left[ (\tilde{\mu} + \mu_t) \left( \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{\partial \tilde{u}_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \frac{\partial \tilde{u}_k}{\partial x_k} \right) - \frac{2}{3} \bar{\rho} \tilde{k} \delta_{ij} \right], \\ \frac{\partial \bar{\rho} \tilde{H}}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{H}}{\partial x_j} &= \frac{\partial \bar{p}}{\partial t} + \frac{\partial}{\partial x_j} \left[ \left( \frac{\tilde{\lambda}}{c_p} + \frac{\mu_t}{Pr_t} \right) \frac{\partial \tilde{h}}{\partial x_j} \right] \\ &\quad + \frac{\partial}{\partial x_j} \left[ (\tilde{\mu} + \sigma_k \mu_t) \frac{\partial \tilde{k}}{\partial x_j} + \tilde{u}_i (\bar{\tau}_{ij} + \bar{\tau}_{ij}^R) \right], \\ \frac{\partial \bar{\rho} \tilde{Y}_k}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{Y}_k}{\partial x_j} &= \frac{\partial}{\partial x_j} \left[ \left( \bar{\rho} \tilde{D}_k + \frac{\tilde{\mu}_t}{Sc_t} \right) \frac{\partial \tilde{Y}_k}{\partial x_j} \right] + \bar{\omega}_k, \\ \frac{\partial \bar{\rho} \tilde{k}}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{k}}{\partial x_j} &= \mathcal{D}_k + P_k^+ - P_k^-, \\ \frac{\partial \bar{\rho} \tilde{\omega}}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{\omega}}{\partial x_j} &= \mathcal{D}_\omega + P_\omega^+ - P_\omega^-, \end{aligned}$$

# Numerical method

- Unsteady RANS using OpenFOAM
- Axisymmetric discretization
- Second-order-accurate Crank-Nicolson time integration
- WENO schemes for advection terms (different orders of accuracy investigated)
- Second-order-accurate linear scheme for diffusion terms
- PIMPLE (combination of PISO and SIMPLE) strategy for pressure-velocity coupling
- rhoReactingPimpleFoam OpenFOAM solver (multi-species solver)

# Numerical method

The main idea of WENO schemes is to employ a polynomial representation of the solution at each cell and compute the average value as:

$$\bar{\Phi}_i = \frac{1}{|V_i|} \int_{V_i} \Phi(\vec{x}) dx dy dz = \frac{1}{|V'_i|} \int_{V'_i} \Phi(\vec{\xi}) d\xi d\eta d\zeta = \frac{1}{|V'_i|} \int_{V'_i} p(\vec{\xi}) d\xi d\eta d\zeta$$

Several polynomial based on different stencils are considered at each cell and a weighted convex combination of them is used.

In the present work the WENO implementation of

## Efficient WENO library for OpenFOAM

Jan Wilhelm Gärtner<sup>a,\*</sup>, Andreas Kronenburg<sup>a</sup>, Tobias Martin<sup>b</sup>

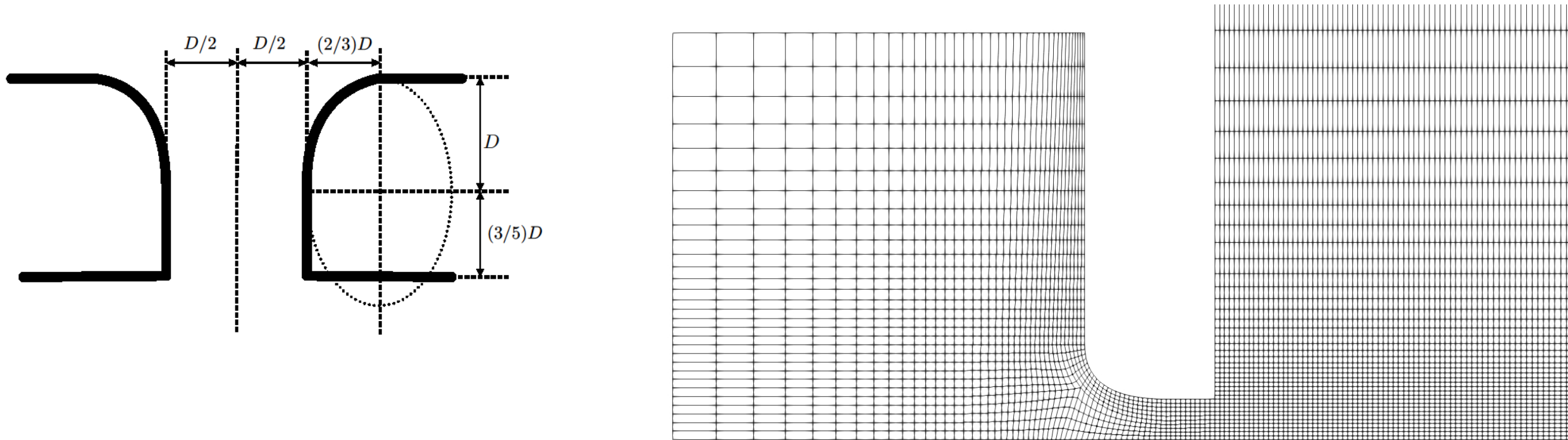
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<sup>b</sup> Department of Civil and Environmental Engineering, Norwegian University of Science and Technology, 7491 Trondheim, Norway

published in **SoftwareX 12 (2020) 100611**, has been employed.

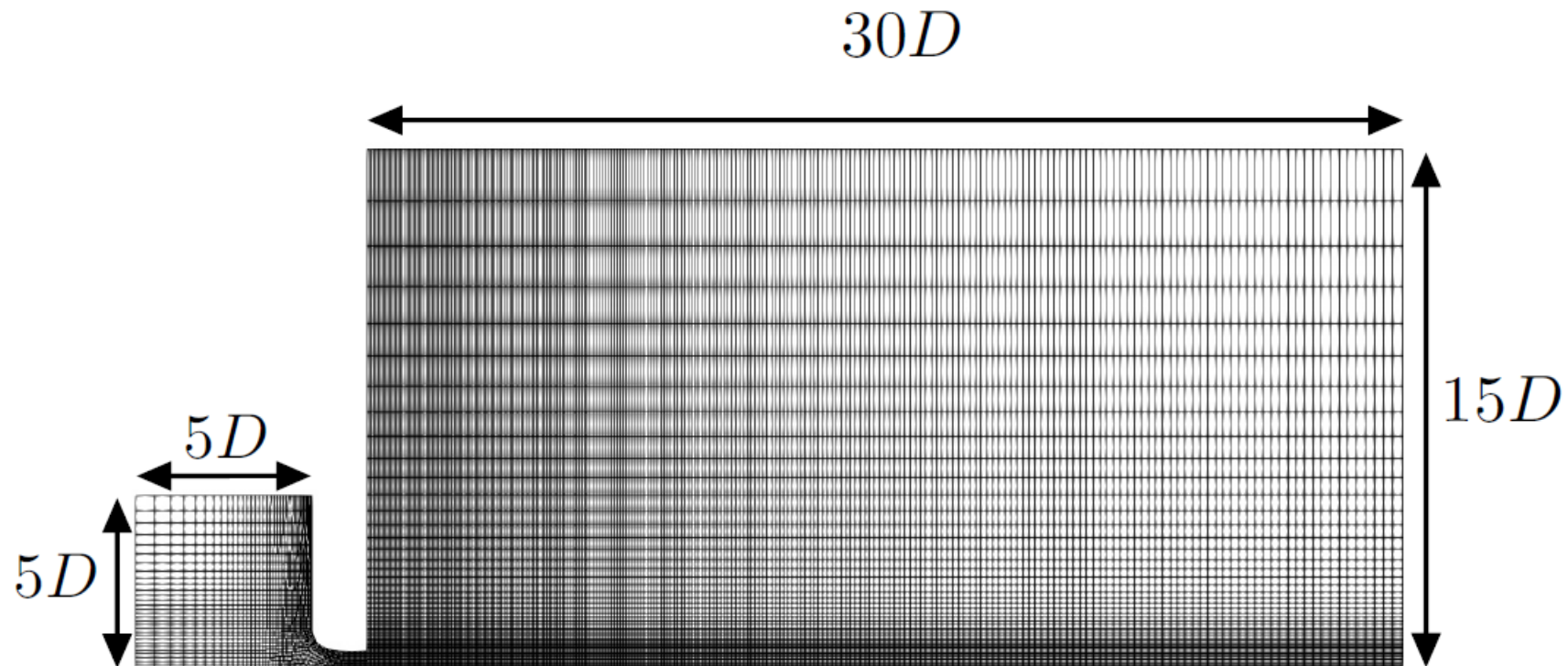
# Simulation set-up

Nozzle geometry ( $D=1.5$  mm, Hamzehloo et al. [4]) and grid



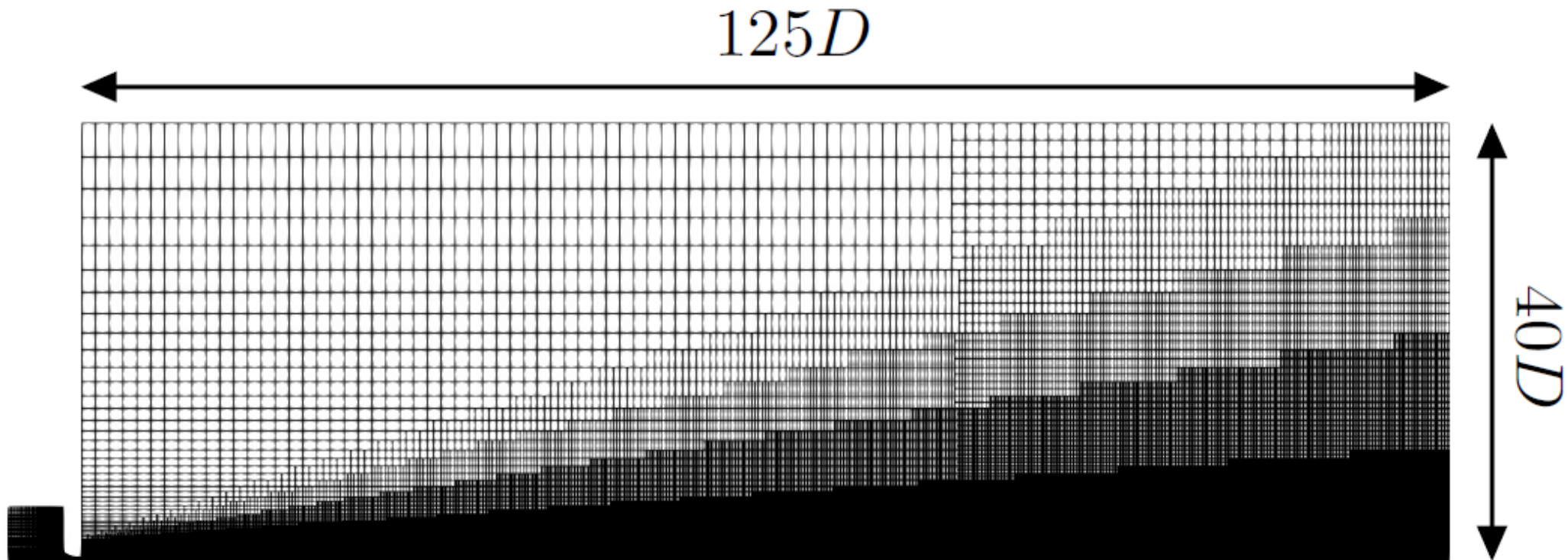
# Simulation set-up

“Short” computational domain dimensions (50905 cells)



# Simulation set-up

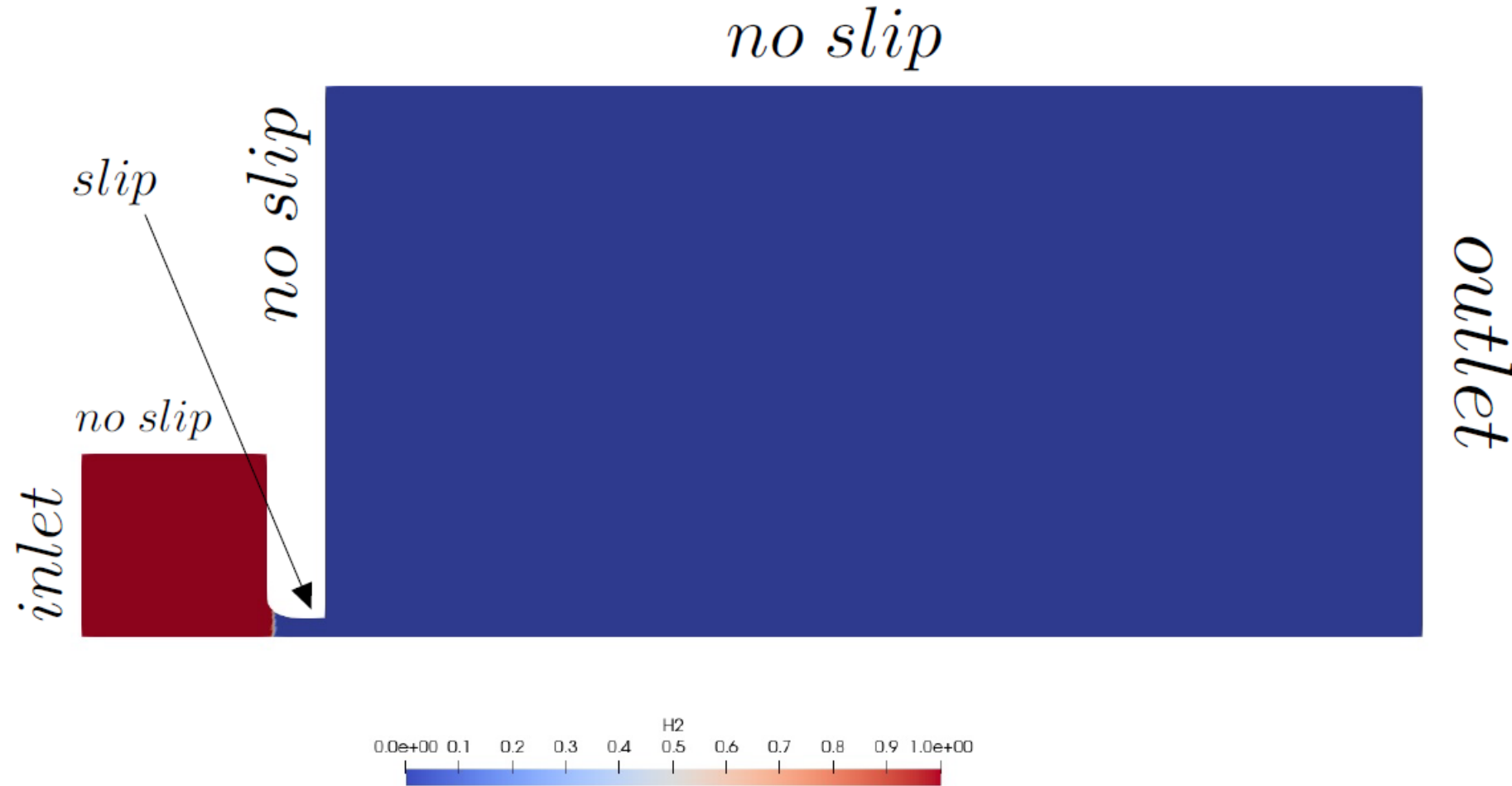
“Long” computational domain dimensions (192512 cells)



# Simulation set-up

## Initial and boundary conditions

- **Initialization** (Vuorinen et al. [2])
- **Nozzle wall:** slip, being  $Kn \approx 0.1 \div 0.01$  (Gad-el-Hak [3], Hamzehloo et al. [4])
- **Inlet:** total pressure, total temperature, flow angle
- **Outlet:** static pressure (98370 Pa)



# Results

Grid refinement study (NPR=10)

Mach disk **height** [mm] for different grids and schemes: experimental value **3.05**  
(Ruggles, A. J. & Ekoto I. W. [5])

|               | <i>D10</i> | <i>D20</i> | <i>D40</i>  | <i>D80</i> |
|---------------|------------|------------|-------------|------------|
| Linear Upwind | 3.62       | 3.62       | 3.23        | 3.07       |
| WENO1         | 3.65       | 3.31       | 3.18        | 3.27       |
| WENO2         | 3.74       | 3.33       | 3.14        | 3.17       |
| WENO3         | 3.79       | 3.28       | <b>3.09</b> | 3.13       |

# Results

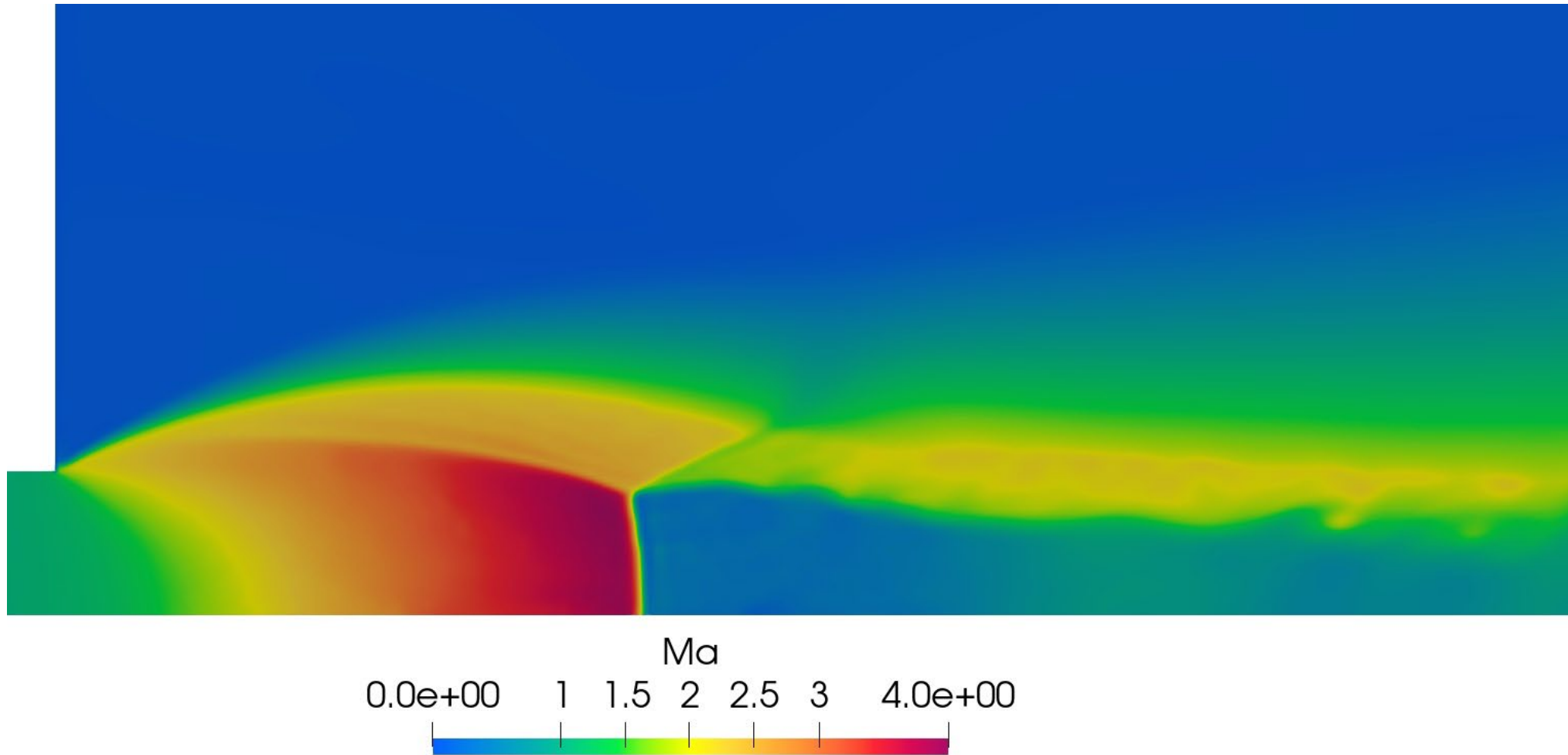
Grid refinement study (NPR=10)

Mach disk **width** [mm] for different grids and schemes: experimental value **1.30**  
(Ruggles, A. J. & Ekoto I. W. [5])

|               | <i>D10</i> | <i>D20</i> | <i>D40</i>  | <i>D80</i> |
|---------------|------------|------------|-------------|------------|
| Linear Upwind | 0.36       | 0.29       | 1.08        | 1.25       |
| WENO1         | 0.39       | 0.89       | 1.25        | 1.36       |
| WENO2         | 0.40       | 0.99       | 1.25        | 1.18       |
| WENO3         | 1.09       | 3.28       | <b>1.27</b> | 1.33       |

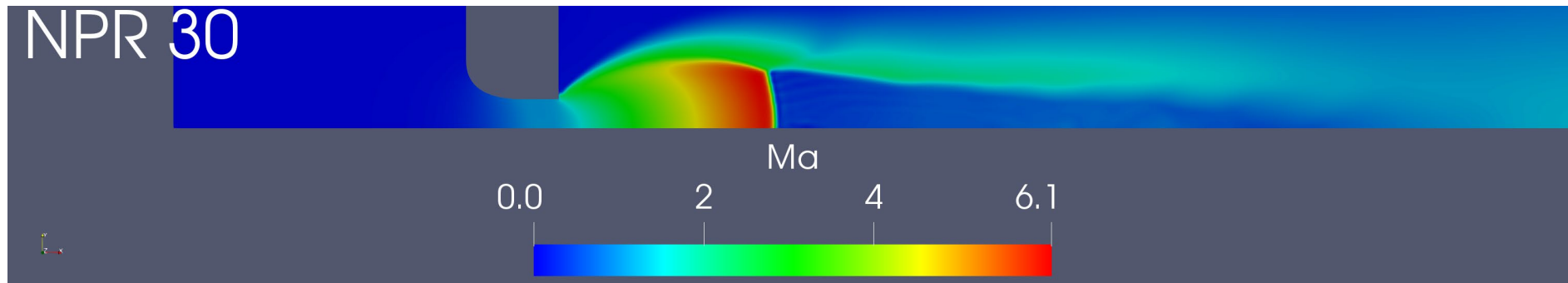
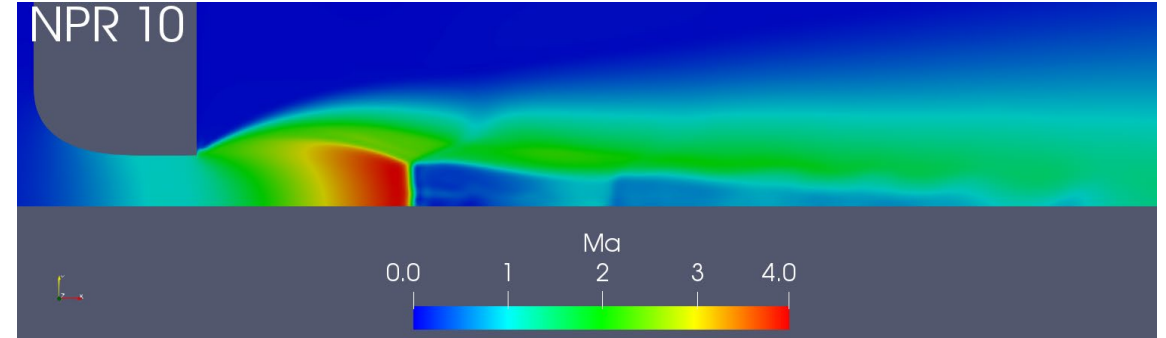
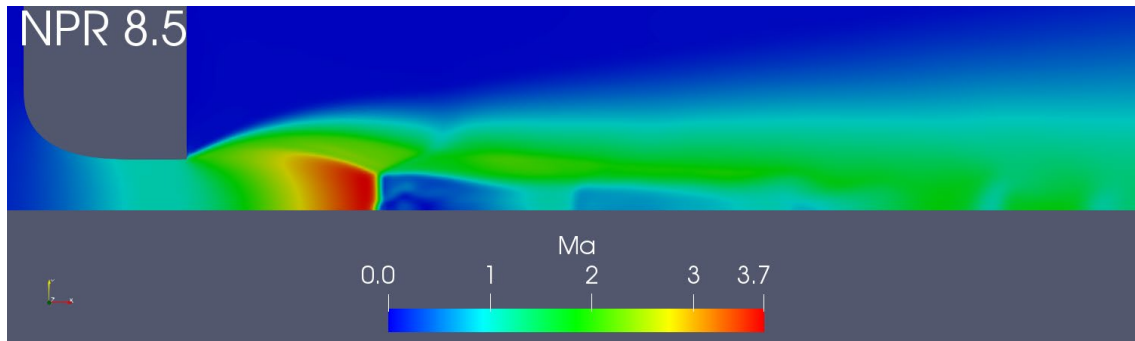
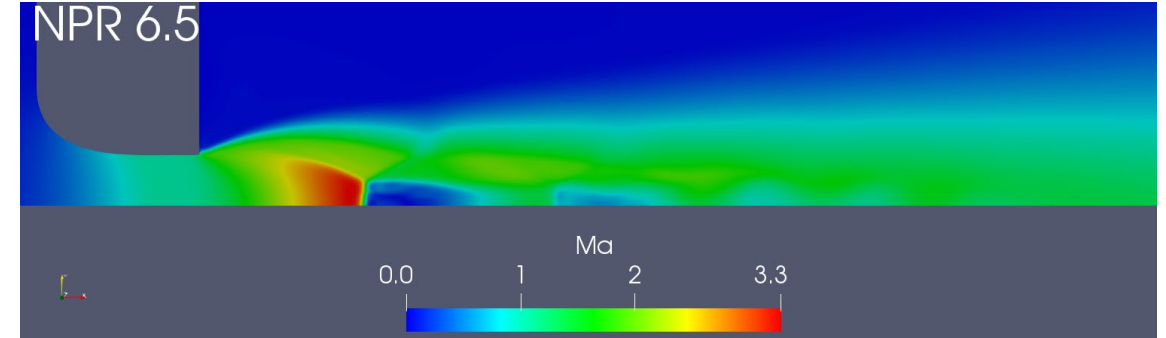
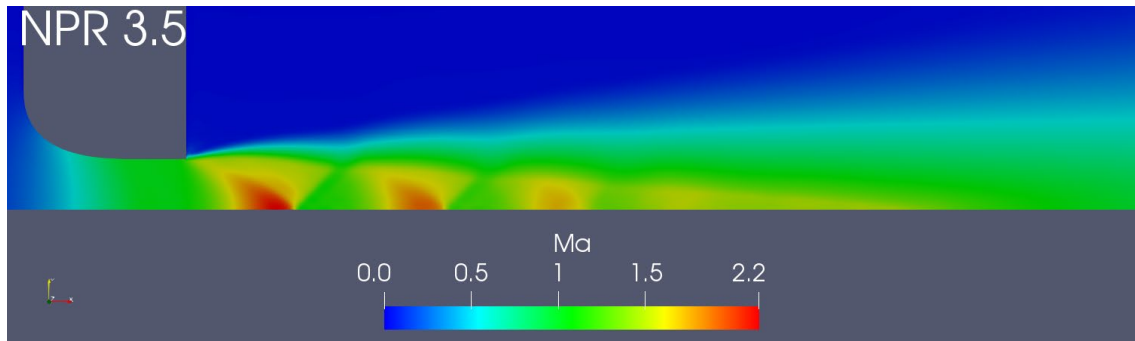
# Results

Barrel shock structure (NPR=10,  $t = 0.2$  [ms], WENO 3-0)

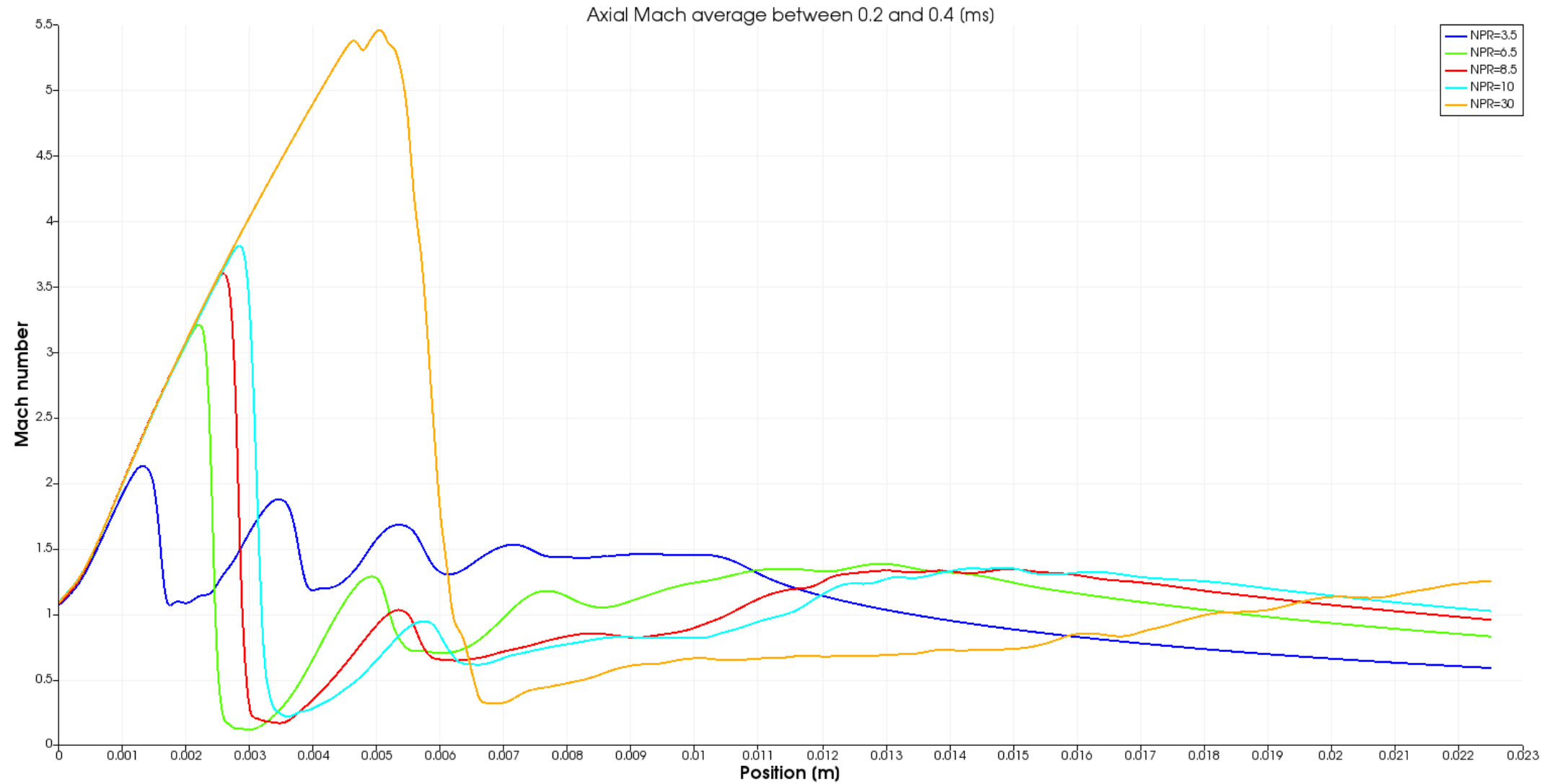


# Results

Mach number ( $t = 0.2$  [ms] WENO 3-0)

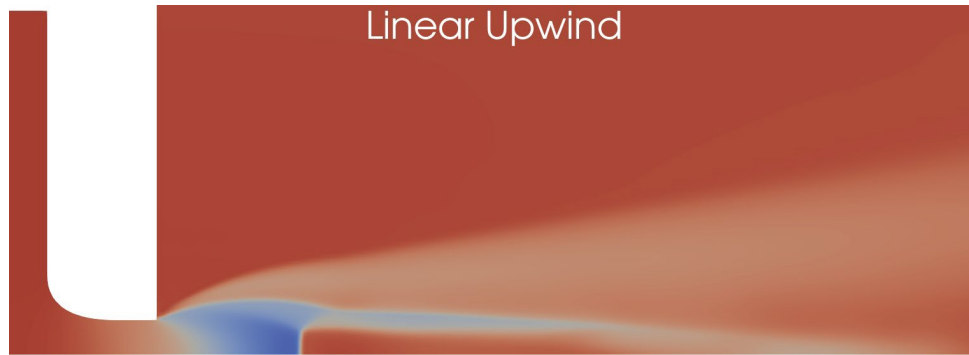


# Results



# Results

Kelvin-Helmholtz instabilities only using WENO3 (Temperature [K], NPR = 10)



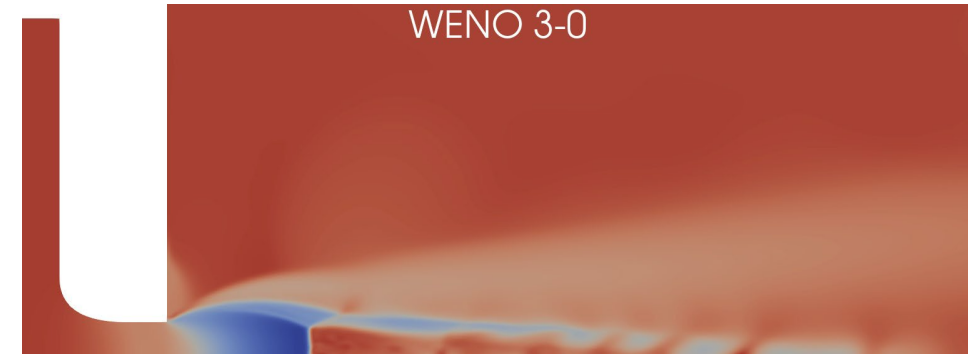
T  
5.7e+01 150 200 250 3.2e+02



T  
5.9e+01 150 200 250 3.2e+02



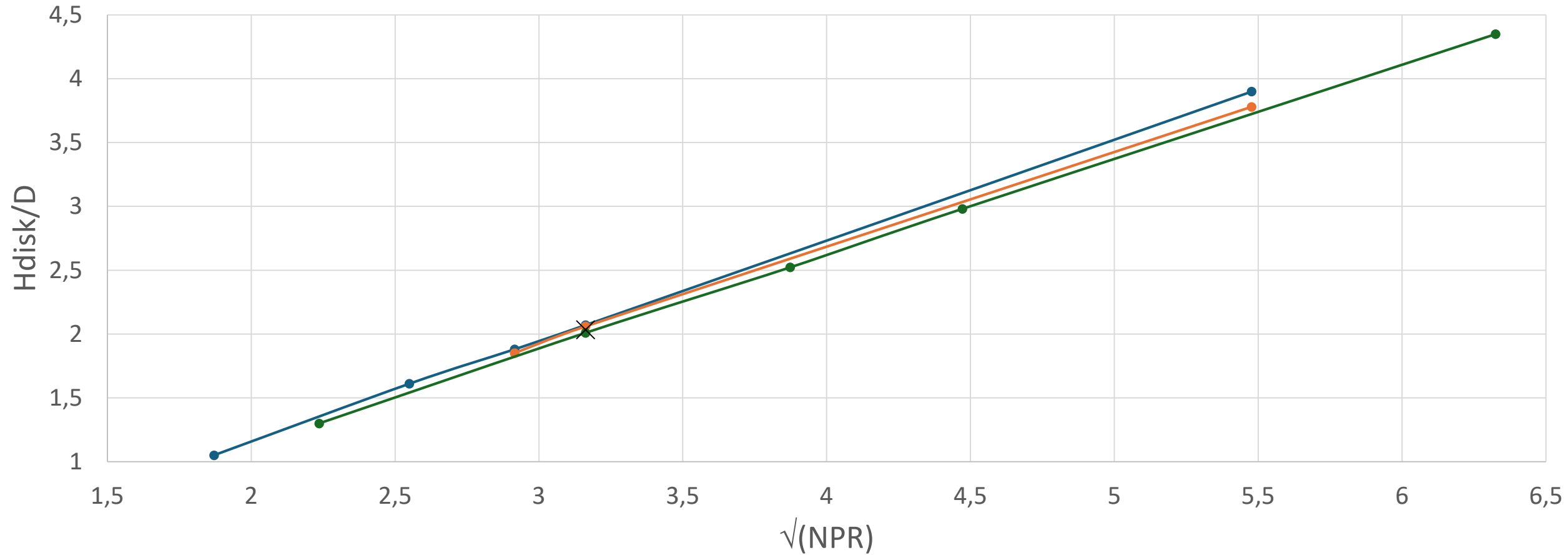
T  
5.7e+01 150 200 250 3.2e+02



T  
5.7e+01 150 200 250 3.2e+02

# Results

Correlation is a straight line



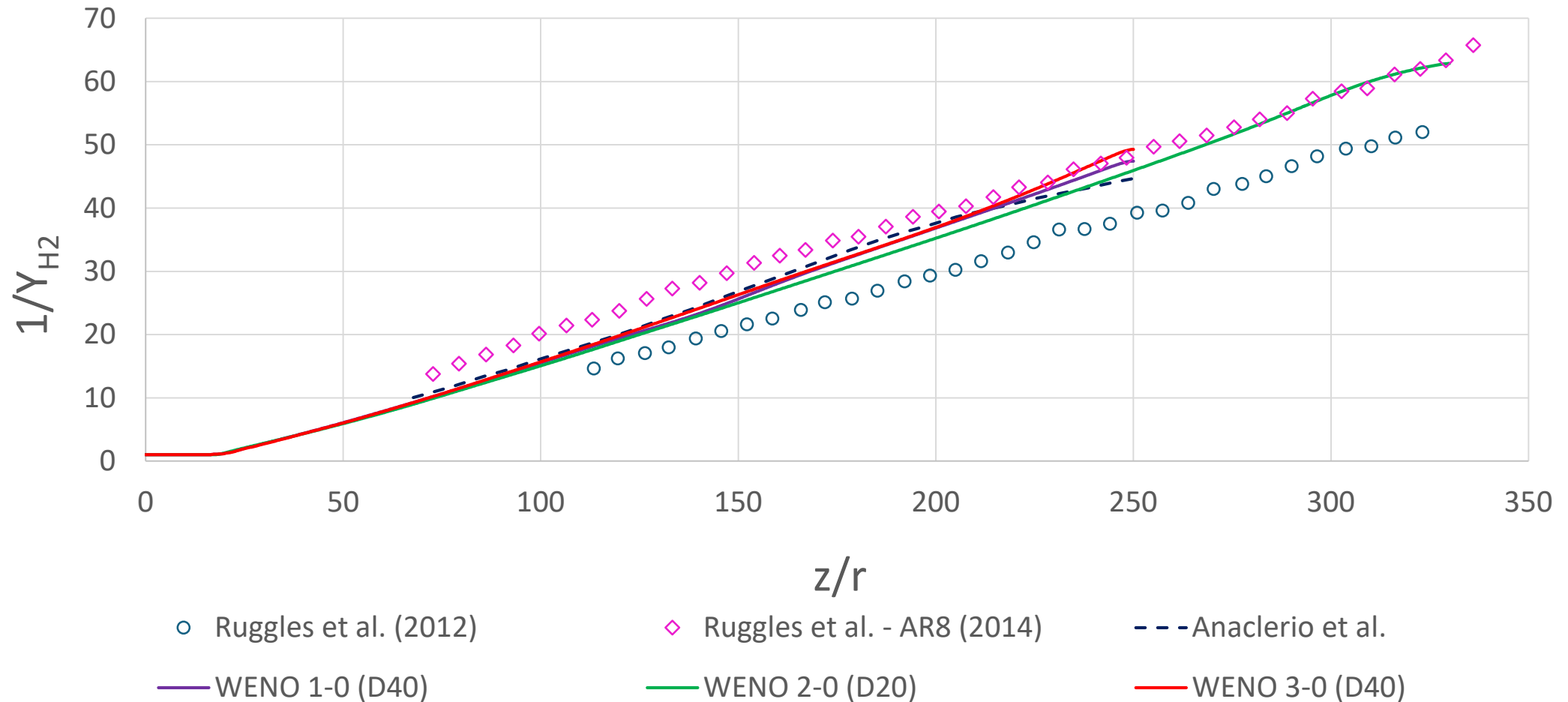
—●— Hdisk/D (URANS WENO 3-0)    —●— Hdisk/D (Hamzehloo et al., [4])    —●— Hdisk/D (Anaclerio et al., [7])    —x— Experimental

# Results

| $NPR$ | <i>Solver</i>              | $T_{min} [K]$ | $U_{max} [m/s]$ | $Ma_{max}$ |
|-------|----------------------------|---------------|-----------------|------------|
| 8.5   | RANS + WENO 3-0            | 69            | 2453            | 3.7        |
|       | LES (Hamzehloo et al. [4]) | 77            | 2493            | 3.72       |
| 10    | RANS + WENO 3-0            | 61            | 2491.6          | 4          |
|       | RANS + WENO 3-1            | 61.6          | 2489.8          | 3.9        |
|       | LES (Hamzehloo et al. [4]) | 71            | 2531            | 3.96       |
| 30    | RANS + WENO 3-0            | 28            | 2637.2          | 6.1        |
|       | RANS + WENO 3-1            | 27.7          | 2637.2          | 6.1        |
|       | LES (Hamzehloo et al. [4]) | 41            | 2695            | 5.53       |

# Results

H<sub>2</sub> mass concentration ( $Y_{H_2}$ ) as a function of distance ( $Z$ ) from nozzle exit (NPR=10)



# Conclusions

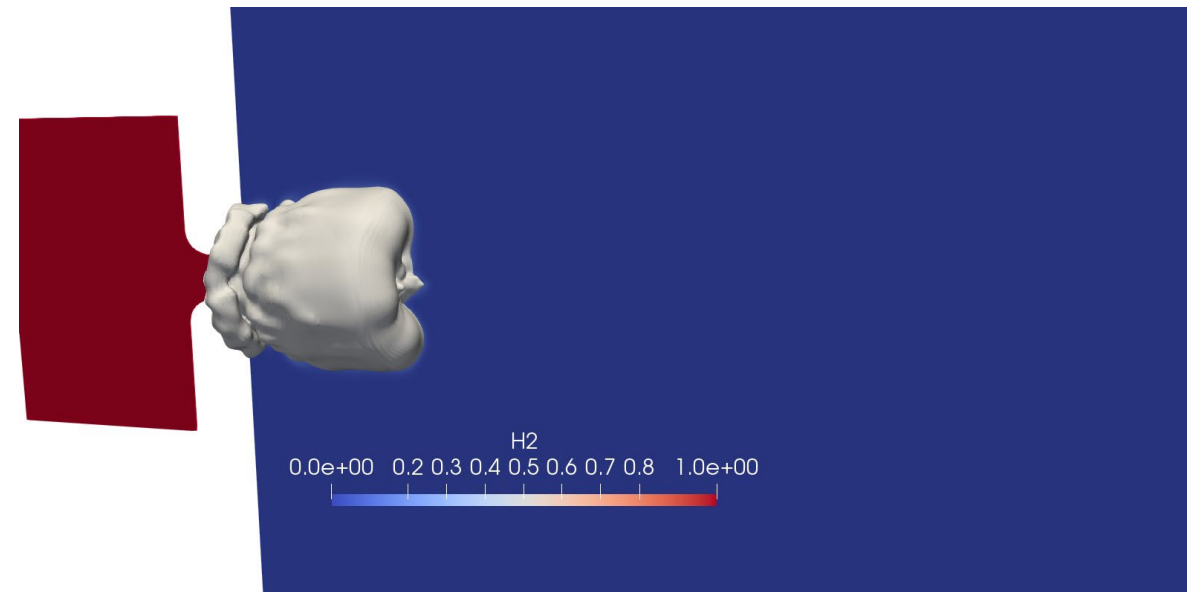
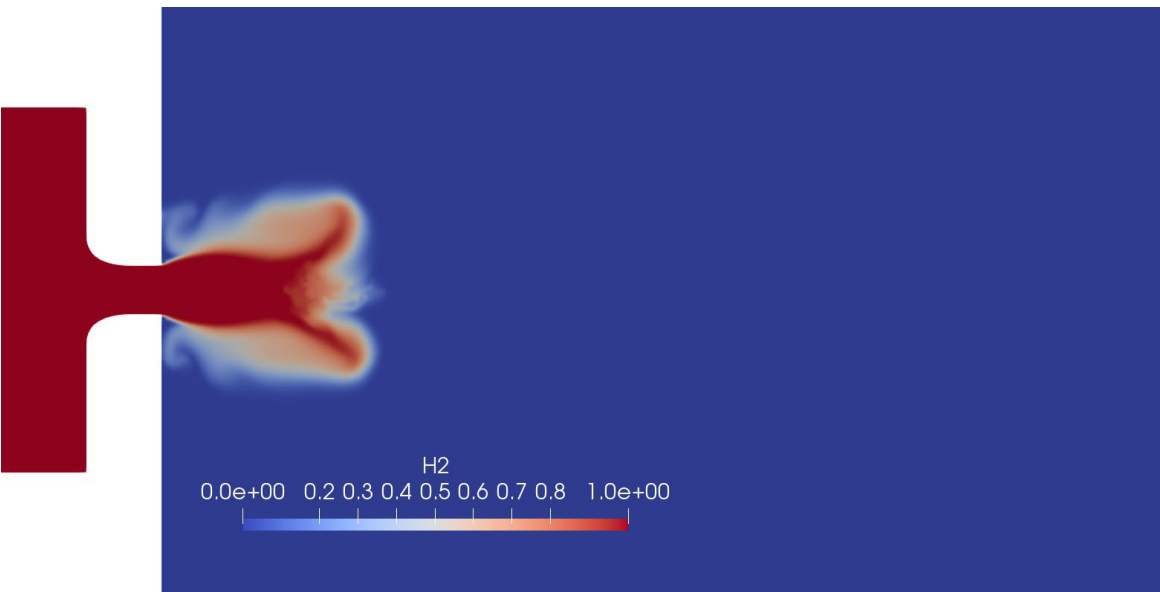
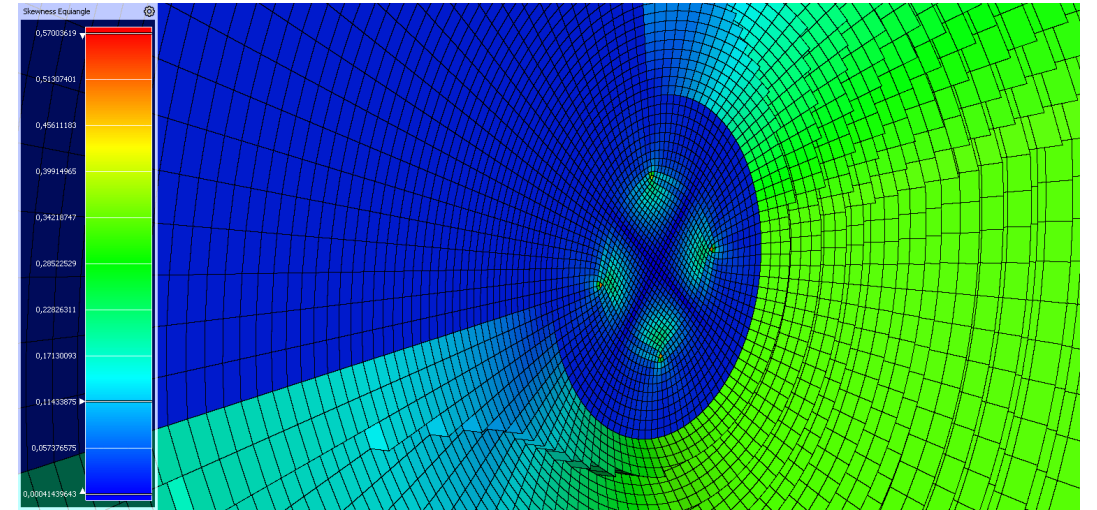
- The near-nozzle under-expanded jet structure and the H<sub>2</sub> jet diffusion have been analyzed by an URANS + WENO model
- The method has been validated by a grid refinement study versus experimental and numerical data.
- Linear correlation for Mach disk height at different NPRs has been verified
- The advantage of employing higher spatial order of accuracy is clearly demonstrated by the capturing of coherent structures
- The model tends to overestimate diffusion

# Future work

- Improve diffusion model (optimize turbulence model coefficient by a Machine Learning technique)
- Extract POD/DMD modes to construct a reduced order model
- Nozzle geometry optimization by adjoint approach with automatic differentiation
- Extension to 3D-URANS
- Extension to LES

# Preliminary 3D Results

~5M cells grid (D40)



# References

- [1] Crist, S., Sherman, P. & Glass, D. Study of the highly underexpanded sonic jet. AIAA Journal 1966 4 (1), 68–71.
- [2] V. Vuorinen, J. Yu, S. Tirunagari, O. Kaario, M. Larmi, C. Duwig, B. J. Boersma. Large-Eddy Simulation of Highly Underexpanded Transient Gas Jets. Physics of Fluids 1 January 2013; 25 (1): 016101
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- [4] Hamzehloo A, Aleiferis PG. *Large eddy simulation of highly turbulent under-expanded hydrogen and methane jets for gaseous-fueled internal combustion engines*. Int. J. Hydrogen Energy, 2014. 39(36):21275-21296.
- [5] Ruggles, A. J. & Ekoto, I. W. 2012 Ignitability and mixing of underexpanded hydrogen jets. International Journal of Hydrogen Energy 37, 17549-17560.
- [6] Ruggles, A. J. & Ekoto, I. W. 2014 Experimental investigation of nozzle aspect ratio effects on underexpanded hydrogen jet release characteristics. International Journal of Hydrogen Energy 39, 20331-20338.
- [7] Anaclerio G, Capurso T, Torresi M, Camporeale SM. *Numerical characterization of hydrogen under-expanded jets with a focus on Internal Combustion Engines applications*. International Journal of Engine Research. 2023; 24(8):3342-3358.